

Anti-Commutator Theorem and Essential Self-Adjointness

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Abstract. In this article, the self-adjoint extensions of symmetric operators satisfying anti-commutation relations are considered. Then it is proven that an anti-commutative type of the Glimm-Jaffe-Nelson commutator theorem follows. In addition we investigate its application to the Dirac operator with external potentials.

1 Introduction and Main Theorem

In this paper we consider the essential self-adjointness of symmetric operators satisfying anti-commutation relations. Let H be a linear operator on a separable Hilbert space \mathcal{H} . It is said that H is symmetric if $H \subset H^*$, that is $\mathcal{D}(H) \subset \mathcal{D}(H^*)$ and $H^*\Psi = H\Psi$ for $\Psi \in \mathcal{D}(H)$, and we say that H is self-adjoint if $H = H^*$. In particular it is said that H is essentially self-adjoint if its closure \overline{H} is self-adjoint. We are interested in conditions under which a symmetric operator has its self-adjoint extension. The Glimm-Jaffe-Nelson commutator theorem ([1]; Theorem 2.32, [7]; Theorem X.36) is one criteria to show the essential self-adjointness for commuting symmetric operators. The commutator theorem shows that if a symmetric operator H and a self-adjoint operator S obey a commutation relation on a dense subspace \mathcal{D} which is a core of S , then H is essentially self-adjoint on \mathcal{D} . Historically, Glimm-Jaffe [3] and Nelson [4] investigated the commutator theorem for the application to quantum field models, and Faris-Lavine [2] considered its application to the quantum mechanical models. We investigate an anti-commutative version of the commutator theorem.

Let us consider symmetric operators X and Y . Then the real part and the imaginary part of the inner product $(X\Psi, Y\Psi)$ for $\Psi \in \mathcal{D}(XY) \cap \mathcal{D}(YX)$ are expressed by

$$\text{Re}(X\Psi, Y\Psi) = \frac{1}{2}(\Psi, \{X, Y\}\Psi), \quad (1)$$

and

$$\text{Im}(X\Psi, Y\Psi) = \frac{1}{2}(\Psi, [X, Y]\Psi), \quad (2)$$

where $\{X, Y\} = XY + YX$ and $[X, Y] = XY - YX$. In the commutator theorem, the imaginary part (2) is estimated. Instead of this, we estimate the real part (1), and obtain an anti-commutative version of the commutator theorem. Here we overview the commutator theorem, and later we state the main theorem.

Let H and O be linear operators on \mathcal{H} . We introduce the following conditions.

(C.1) H is symmetric and O is self-adjoint.

(C.2) There exists $\delta_0 > 0$ such that for $\Psi \in \mathcal{D}(O)$,

$$\delta_0(\Psi, \Psi) \leq |(\Psi, O\Psi)|.$$

(C.3) O has a core \mathcal{D}_0 satisfying $\mathcal{D}_0 \subset \mathcal{D}(H)$, and there exist constants $a \geq 0$ and $b \geq 0$ such that for all $\Psi \in \mathcal{D}_0$,

$$\|H\Psi\| \leq a\|O\Psi\| + b\|\Psi\|.$$

Proposition A (Glimm-Jaffe-Nelson Commutator Theorem)

Let H and O be the operators satisfying (C.1)-(C.3). In addition we suppose (i) or (ii) below :

(i) There exists a constant $c_+ > 0$ such that for all $\Psi \in \mathcal{D}_0$,

$$\left| (H\Psi, O\Psi) - (O\Psi, H\Psi) \right| \leq c_+ |(\Psi, O\Psi)|. \quad (3)$$

(ii) There exists a constant $c_- > 0$ such that for all $\Psi \in \mathcal{D}_0$,

$$c_- |(\Psi, O\Psi)| \leq \left| (H\Psi, O\Psi) - (O\Psi, H\Psi) \right|. \quad (4)$$

Then H is essentially self-adjoint on \mathcal{D}_0 .

Remark 1 In the commutator theorem, the condition (i) is usually supposed. But we will see that the commutator theorem also follows under the condition (ii).

(Proof of Proposition A) Since the proof with the condition (i) is well-known, let us only give the proof by supposing (ii). From the general theorem ([7], Theorem X.1), it is enough to show that for some $z \in \mathbf{C} \setminus \mathbf{R}$, $\dim \ker((H_{|\mathcal{D}_0})^* + z^\sharp) = 0$ where $z^\sharp = z, z^*$. Since \mathcal{D}_0 is a core of O , it is seen from (C.3) and (ii) that $\mathcal{D}(O) \subset \mathcal{D}(\overline{H_{|\mathcal{D}_0}})$ and for $\Phi \in \mathcal{D}(O)$,

$$c_- |(\Phi, O\Phi)| \leq \left| (\overline{H_{|\mathcal{D}_0}}\Phi, O\Phi) - (O\Phi, \overline{H_{|\mathcal{D}_0}}\Phi) \right|. \quad (5)$$

Let $\Psi \in \mathcal{D}((H_{|\mathcal{D}_0})^*)$ and $\Xi = O^{-1}\Psi \in \mathcal{D}(O)$. Here it follows from (C.2) that O^{-1} is bounded. By using $((H_{|\mathcal{D}_0})^*)^* = \overline{H_{|\mathcal{D}_0}}$, we obtain that

$$\text{Im} \left(\Xi, ((H_{|\mathcal{D}_0})^* + z^\sharp)\Psi \right) = \frac{1}{2} \left((\overline{H_{|\mathcal{D}_0}}\Xi, O\Xi) - (O\Xi, \overline{H_{|\mathcal{D}_0}}\Xi) \right) + \text{Im} z^\sharp (\Xi, O\Xi). \quad (6)$$

Then for $z \in \mathbf{C} \setminus \mathbf{R}$ satisfying $|\text{Im} z| < \frac{c_-}{2}$, we have from (5) and (6), that

$$\left| \text{Im} \left(\Xi, ((H_{|\mathcal{D}_0})^* + z^\sharp)\Psi \right) \right| \geq \left(\frac{c_-}{2} - |\text{Im} z| \right) |(\Xi, O\Xi)| \geq \delta_O \left(\frac{c_-}{2} - |\text{Im} z| \right) (\Xi, \Xi). \quad (7)$$

Hence, if $\Psi \in \ker((H_{|\mathcal{D}_0})^* + z^\sharp)$, we have $\Xi = O^{-1}\Psi = 0$ from (7), and then $\Psi = 0$. Thus the proof is obtained. ■

Here we state the main theorem.

Theorem 1 Assume **(C.1)-(C.3)**. In addition we suppose the following **(I)** or **(II)** holds.

(I) There exists a constant $d_+ > 0$ such that for all $\Psi \in \mathcal{D}_0$,

$$\left| (H\Psi, O\Psi) + (O\Psi, H\Psi) \right| \leq d_+ |(\Psi, O\Psi)|. \quad (8)$$

(II) There exists a constant $d_- > 0$ such that for all $\Psi \in \mathcal{D}_0$,

$$d_- |(\Psi, O\Psi)| \leq \left| (H\Psi, O\Psi) + (O\Psi, H\Psi) \right|. \quad (9)$$

Then H is essentially self-adjoint on \mathcal{D}_0 .

(Proof of Theorem 1) In a similar way to Proposition A, let us show that for some $z \in \mathbf{C} \setminus \mathbf{R}$, $\dim \ker((H_{|\mathcal{D}_0})^* + z^\sharp) = 0$ where $z^\sharp = z, z^*$. Let $\Psi \in \mathcal{D}((H_{|\mathcal{D}_0})^*)$ and $\Xi = O^{-1}\Psi \in \mathcal{D}(O)$. Then by $((H_{|\mathcal{D}_0})^*)^* = \overline{H_{|\mathcal{D}_0}}$, we see that

$$\operatorname{Re} \left(\Xi, ((H_{|\mathcal{D}_0})^* + z^\sharp)\Psi \right) = \frac{1}{2} \left((\overline{H_{|\mathcal{D}_0}}\Xi, A\Xi) + (O\Xi, \overline{H_{|\mathcal{D}_0}}\Xi) \right) + \operatorname{Re} z (\Xi, O\Xi). \quad (10)$$

First we assume that **(I)** holds. Let us take $z \in \mathbf{C} \setminus \mathbf{R}$ satisfying $|\operatorname{Re} z| > \frac{d_+}{2}$. Since \mathcal{D}_0 is a core of O , it follows from **(C.3)** and **(I)** that $\mathcal{D}(O) \subset \mathcal{D}(\overline{H_{|\mathcal{D}_0}})$ and for $\Phi \in \mathcal{D}(O)$,

$$\left| (\overline{H_{|\mathcal{D}_0}}\Phi, O\Phi) + (O\Phi, \overline{H_{|\mathcal{D}_0}}\Phi) \right| \leq d_+ |(\Phi, O\Phi)|. \quad (11)$$

By (10) and (11), we have

$$\left| \operatorname{Re} \left(\Xi, ((H_{|\mathcal{D}_0})^* + z^\sharp)\Psi \right) \right| \geq \left(|\operatorname{Re} z| - \frac{d_+}{2} \right) |(\Xi, O\Xi)| \geq \delta_0 \left(|\operatorname{Re} z| - \frac{d_+}{2} \right) (\Xi, \Xi). \quad (12)$$

Hence if $\Psi \in \ker((H_{|\mathcal{D}_0})^* + z^\sharp)$, we have $\Xi = O^{-1}\Psi = 0$ from (12). Then we have $\Psi = 0$. Next we suppose that **(II)** follows. Let us take $z \in \mathbf{C} \setminus \mathbf{R}$ satisfying $|\operatorname{Re} z| < \frac{d_-}{2}$. Since \mathcal{D}_0 is a core of O , it also follows from **(C.3)** and **(II)** that $\mathcal{D}(O) \subset \mathcal{D}(\overline{H_{|\mathcal{D}_0}})$ and for $\Phi \in \mathcal{D}(O)$,

$$d_- |(\Phi, O\Phi)| \leq \left| (\overline{H_{|\mathcal{D}_0}}\Phi, O\Phi) + (O\Phi, \overline{H_{|\mathcal{D}_0}}\Phi) \right|. \quad (13)$$

Then from (10) and (13), we have

$$\left| \operatorname{Re} \left(\Xi, ((H_{|\mathcal{D}_0})^* + z^\sharp)\Psi \right) \right| \geq \left(\frac{d_-}{2} - |\operatorname{Re} z| \right) |(\Xi, O\Xi)| \geq \delta_0 \left(\frac{d_-}{2} - |\operatorname{Re} z| \right) (\Xi, \Xi). \quad (14)$$

Hence, if $\Psi \in \ker((H_{|\mathcal{D}_0})^* + z^\sharp)$, we have $\Psi = 0$ since $\Xi = O^{-1}\Psi = 0$ from (14). Thus the proof is obtained. ■

Here we consider the following condition **(C.2)'**, which is weaker than **(C.2)**.

(C.2)' There exist $e > 0$ such that for $\Phi \in \mathcal{D}(O)$,

$$e\|\Phi\| \leq \|O\Phi\|.$$

Then the following Proposition 2 holds. It is noted that the conditions **(A)** and **(B)** are stronger than **(ii)** in Proposition A and **(II)** in Theorem 1, respectively.

Proposition 2

Assume **(C.1), (C.2)', and (C.3)**. In addition the following **(A)** or **(B)** holds.

(A) There exists a constant $\mu > 0$ such that for all $\Psi \in \mathcal{D}_0$,

$$\mu \|O\Psi\|^2 \leq |(H\Psi, O\Psi) - (O\Psi, H\Psi)|. \quad (15)$$

(B) There exists a constant $v > 0$ such that for all $\Psi \in \mathcal{D}_0$,

$$v \|O\Psi\|^2 \leq |(H\Psi, O\Psi) + (O\Psi, H\Psi)|. \quad (16)$$

Then H is essentially self-adjoint on \mathcal{D}_0 .

(Proof of Proposition 2) Let us prove that $\dim \ker((H_{|\mathcal{D}_0})^* + z^\sharp) = 0$ for some $z \in \mathbf{C} \setminus \mathbf{R}$. Let $\Psi \in \mathcal{D}((H_{|\mathcal{D}_0})^*)$ and $\Xi = O^{-1}\Psi \in \mathcal{D}(O)$. First we suppose **(A)**. Since \mathcal{D}_0 is a core of O , it is seen from **(C.3)** and **(A)** that $\mathcal{D}(O) \subset \mathcal{D}(\overline{H_{|\mathcal{D}_0}})$ and for $\Phi \in \mathcal{D}(O)$,

$$\mu \|O\Phi\|^2 \leq |(\overline{H_{|\mathcal{D}_0}}\Phi, O\Phi) - (O\Phi, \overline{H_{|\mathcal{D}_0}}\Phi)|. \quad (17)$$

Let $z \in \mathbf{C} \setminus \mathbf{R}$ satisfying $|\operatorname{Im} z| < \frac{v_e}{2}$. Then, by using $((H_{|\mathcal{D}_0})^*)^* = \overline{H_{|\mathcal{D}_0}}$, we have

$$|\operatorname{Im}(\Xi, ((H_{|\mathcal{D}_0})^* + z^\sharp)\Psi)| \geq |\operatorname{Im}(\overline{H_{|\mathcal{D}_0}}\Xi, O\Xi)| - |\operatorname{Im} z| |(\Xi, O\xi)| \geq \frac{\mu}{2} \|O\Xi\|^2 - |\operatorname{Im} z| |(\Xi, O\xi)|. \quad (18)$$

Note that $|(\Xi, O\Xi)| \leq \|\Xi\| \|O\Xi\| \leq \frac{1}{e} \|O\Xi\|^2$ from **(C.2)'**. Then by (18) we have

$$|\operatorname{Im}(\Xi, ((H_{|\mathcal{D}_0})^* + z^\sharp)\Psi)| \geq \left(\frac{v}{2} - \frac{|\operatorname{Im} z|}{e} \right) \|O\Xi\|^2 = \left(\frac{v}{2} - \frac{|\operatorname{Im} z|}{e} \right) \|\Psi\|^2. \quad (19)$$

Then if $\Psi \in \ker((H_{|\mathcal{D}_0})^* + z^\sharp)$, we have $\Psi = 0$ from (19). Next we suppose **(B)**. Since \mathcal{D}_0 is a core of O , it is also seen from **(C.3)** and **(B)** that $\mathcal{D}(O) \subset \mathcal{D}(\overline{H_{|\mathcal{D}_0}})$ and for $\Phi \in \mathcal{D}(O)$,

$$v \|O\Phi\|^2 \leq |(\overline{H_{|\mathcal{D}_0}}\Phi, O\Phi) + (O\Phi, \overline{H_{|\mathcal{D}_0}}\Phi)| \quad (20)$$

Let $z \in \mathbf{C} \setminus \mathbf{R}$ satisfying $|\operatorname{Re} z| < \frac{v_e}{2}$. Then, we see that

$$|\operatorname{Re}(\Xi, ((H_{|\mathcal{D}_0})^* + z^\sharp)\Psi)| \geq |\operatorname{Re}(\overline{H_{|\mathcal{D}_0}}\Xi, O\xi)| - |\operatorname{Re} z| |(\Xi, O\xi)| \geq \frac{v}{2} \|O\Xi\|^2 - |\operatorname{Re} z| |(\Xi, O\xi)|. \quad (21)$$

Since $|(\Xi, O\xi)| \leq \|\Xi\| \|O\Xi\| \leq \frac{1}{e} \|O\Xi\|^2$, we have from (21) that

$$\left| \operatorname{Re}(\Xi, ((H_{|\mathcal{D}_0})^* + z^\sharp)\Psi) \right| \geq \left(\frac{\nu}{2} - \frac{|\operatorname{Re}z|}{e} \right) \|O\Xi\|^2 = \left(\frac{\nu}{2} - \frac{|\operatorname{Re}z|}{e} \right) \|\Psi\|^2. \quad (22)$$

Then if $\Psi \in \ker((H_{|\mathcal{D}_0})^* + z^\sharp)$, we have $\Psi = 0$ from (22). Thus the proof is obtained. ■

2 Application

Let us consider the essential self-adjointness of Dirac operators with external potentials. Let us set the state space by $\mathcal{H}_D = \mathbf{C}^4 \otimes L^2(\mathbf{R}^3)$, and the free Hamiltonian by

$$H_0 = \sum_{j=1}^3 \alpha^j \otimes \hat{P}_j + \beta \otimes M, \quad M > 0,$$

where $\hat{P}_j = -i\partial_{x^j}$, and α^j , $j = 1, 2, 3$, β are 4×4 hermitian matrices satisfying $\{\alpha^j, \alpha^l\} = 2\delta_{j,l}$, $\{\alpha^j, \beta\} = 0$, and $\beta^2 = I$. It is seen that H_0 is self-adjoint on $\mathcal{D}(H_0) = \cap_{j=1}^3 \mathcal{D}(\hat{P}_j)$, and essentially self-adjoint on $\mathcal{D}_c = \mathbf{C}^4 \otimes C_0^\infty(\mathbf{R}^3)$. From the functional calculus we see that

$$M\|\Psi\| \leq \|H_0\Psi\|, \quad \Psi \in \mathcal{D}(H_0). \quad (23)$$

Let us set the total Hamiltonian by

$$H(\kappa) = H_0 + \kappa V, \quad \kappa \in \mathbf{R},$$

where V is a symmetric operator satisfying the condition below :

(D.1) $\mathcal{D}_c \subset \mathcal{D}(V)$, and there exist $a > 0$ and $b \geq 0$ such that for $\Psi \in \mathcal{D}_c$,

$$\|V\Psi\| \leq a\|H_0\Psi\| + b\|\Psi\|.$$

Then, the Kato-Rellich theorem shows that for $|\kappa| < \frac{1}{a}$, $H(\kappa)$ is self-adjoint on $\mathcal{D}(H_0)$, and essentially self-adjoint on \mathcal{D}_c . Here we consider another proof by applying Proposition 2. We assume the additional condition :

(D.2) There exists $c > 0$ such that for $\Psi \in \mathcal{D}_c$,

$$\left| (H_0\Psi, V\Psi) + (V\Psi, H_0\Psi) \right| \leq c\|H_0\Psi\|^2.$$

Then we obtain the corollary below :

Corollary 3

Assume **(D.1)** and **(D.2)**. Then for $|\kappa| < \frac{2}{c}$, $H(\kappa)$ is essentially self-adjoint on \mathcal{D}_c .

(Proof) Let us apply $H(\kappa)$ to H and H_D to O in Proposition 2. We see that **(C.1)** follows, and **(C.2)'** holds from (23). By **(D.1)** and (23), we have $\|H(\kappa)\Psi\| \leq (1 + |\kappa|a) \|H_0\Psi\| + b\|\Psi\|$ for $\Psi \in \mathcal{D}_c$. Thus **(C.3)** follows. We also see from **(D.2)** that for $\Psi \in \mathcal{D}_c$,

$$|(H(\kappa)\Psi, H_0\Psi) + (H_0\Psi, H(\kappa)\Psi)| \geq 2\|H_0\Psi\|^2 - |\kappa| \left((V\Psi, H_0\Psi) + (H_0\Psi, V\Psi) \right) \geq (2 - |\kappa|c)\|H_0\Psi\|^2.$$

Then the condition **(B)** in Proposition 2 is satisfied. Thus we obtain the corollary. ■

Let us consider an example of V in Corollary 3. Let us set

$$V = \tau \otimes K$$

where $\tau = \beta\alpha^1\alpha^2\alpha^3$, and K is the symmetric operator satisfying the conditions below :

(E.1) $C_0^\infty(\mathbf{R}^3) \subset \mathcal{D}(K)$, and there exist $a_1 \geq 0$ and $b_1 \geq 0$ such that

$$\|K\Psi\|^2 \leq a_1 \sum_{j=1}^3 \|\hat{P}_j\Psi\|^2 + b_1\|\Psi\|^2, \quad \Psi \in C_0^\infty(\mathbf{R}^3).$$

(E.2) $K(C_0^\infty(\mathbf{R}^3)) \subset C_0^\infty(\mathbf{R}^3)$, and there exist $a_2 \geq 0$ and $b_2 \geq 0$ such that

$$\sum_{j=1}^3 \| [K, \hat{P}_j] \Psi \|^2 \leq a_2 \sum_{j=1}^3 \|\hat{P}_j\Psi\|^2 + b_2\|\Psi\|^2, \quad \Psi \in C_0^\infty(\mathbf{R}^3).$$

For an example of K , we take $K = P(\hat{P}_1, \hat{P}_2, \hat{P}_3)$ where P is a polynomial of order one. Now let us show that V satisfies **(D.1)**-**(D.2)**. Note that τ is hermitian and $\tau^2 = I$. Then by **(E.1)**, we have

$$\|H(\kappa)\Psi\| \leq \|H_0\Psi\| + |\kappa| \|I \otimes K\Psi\| \leq \left(1 + |\kappa|a_1^{1/2}\right) \|H_0\Psi\| + b_1^{1/2}\|\Psi\|, \quad (24)$$

for $\Psi \in \mathcal{D}_c$. Thus **(D.1)** follows. Note that $\{\tau, \beta\} = 0$ and $\{\tau, \alpha^j\} = 0$, $j = 1, 2, 3$. Then by using $\{X_1 \otimes Y_1, X_2 \otimes Y_2\} = \{X_1, X_2\} \otimes Y_1Y_2 + X_2X_1 \otimes [Y_2, Y_1]$, we have

$$\left| ((\tau \otimes K)\Psi, H_0\Psi) + (H_0\Psi, (\tau \otimes K)\Psi) \right| = \left| \sum_{j=1}^3 (\Psi, ((\alpha^j \tau) \otimes [\hat{P}_j, V_0])\Psi) \right| \leq \|\Psi\| \sum_{j=1}^3 \| [V_0, \hat{P}_j] \Psi \| \quad (25)$$

Then by (23) and **(E.2)**, we have

$$\left| ((\tau \otimes K)\Psi, H_0\Psi) + (H_0\Psi, (\tau \otimes K)\Psi) \right| \leq \frac{1}{M} \left(a_2 + \frac{b_2}{M^2} \right)^{1/2} \|H_D\Psi\|^2. \quad (26)$$

Then from corollary 3, for $|\kappa| < \frac{2M}{\sqrt{a_2 + \frac{b_2}{M^2}}}$, $H_D(\kappa)$ is essentially self-adjoint on D_c .

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